

Fluorescent Derivatives of the GFP Chromophore Give a New Insight into the GFP Fluorescence Process

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ABSTRACT The photophysical properties of synthetic compounds derived from the imidazolidinone chromophore of the green fluorescent protein were determined. Various electron-withdrawing or electron-donating substituents were introduced to mimic the effect of the chromophore surroundings in the protein. The absorption and emission spectra as well as the fluorescence quantum yields in dioxane and glycerol were shown to be highly dependent on the electronic properties of the substituents. We propose a kinetic scheme that takes into account the temperature-dependent twisting of the excited molecule. If the activation energy is low, the molecule most often undergoes an excited-state intramolecular twisting that leads it to the ground state through an avoided crossing between the S_1 and S_0 energy surfaces. For a high activation energy, the torsional motion within the compounds is limited and the ground-state recovery will occur preferentially by fluorescence emission. The excellent correlation between the fluorescence quantum yields and the calculated activation energies to torsion points to the above-mentioned avoided crossing as the main nonradiative deactivation channel in these compounds. Finally, our results are discussed with regard to the chromophore in green fluorescent protein and some of its mutants.

INTRODUCTION

The green fluorescent protein (GFP) from the jellyfish *Aequorea victoria* has become a widely used marker in molecular and cell biology due to its strong intrinsic visible fluorescence, which is easily detectable by fluorescence spectroscopy (for reviews, see Cubitt et al., 1995; Tsien, 1998; Palm and Wlodawer, 1999; Prendergast, 1999; Zimmer, 2002). This protein of around 27 kDa can be fused to many other proteins, allowing their visualization in living cells without interfering with their function. Therefore, it becomes possible to follow signaling and trafficking in cells (Chalfie et al., 1994; Prasher, 1995) and to study protein-protein interactions (Mitra et al., 1996; Park and Raines, 1997; Ozawa et al., 2000). Some GFP variants are also used as noninvasive intracellular pH biosensors (Kneen et al., 1998; Llopis et al., 1998; Robey et al., 1998; Hanson et al., 2002) or fluorescent indicators for local Ca^{2+} concentrations (Miyawaki et al., 1997; Romoser et al., 1997).

The chromophore of this soluble globular protein of 238 amino acid residues is a *p*-hydroxybenzylidene-imidazolidinone derivative formed by an autocatalytic, posttranslational cyclization and oxidation of the polypeptide backbone, involving Ser-65, Tyr-66, and Gly-67 residues (Shimomura, 1979; Cody et al., 1993; Heim et al., 1994). The GFP crystal structure as a monomer (Brejc et al., 1997), as a dimer (Yang et al., 1996), and of several mutants (Ormö et al., 1996; Palm et al., 1997; Wachter et al., 1997) show that the chromophore is always located in the middle part of a central helix inside an 11-stranded β -barrel. It is thus totally embedded in the

protein matrix and isolated from the bulk solvent. The chromophore of the wild type GFP (wt-GFP) possesses peculiar spectroscopic properties: two absorption peaks at 395 and 475 nm, usually attributed respectively to the neutral and anionic forms of the chromophore (Warren and Zimmer, 2001), and one emission peak at 508 nm, with a high fluorescence quantum yield ($\phi = 0.79$) (Ward and Bokman, 1982; Niwa et al., 1996). The substitution of one or more amino acids within the chromophore or in its immediate proximity has allowed the development of GFP mutants with specific absorption and emission properties (Palm et al., 1997; Wachter et al., 1998; Ito et al., 1999).

The photophysical parameters that define the GFP spectroscopic properties are currently under investigation. Vibrational spectroscopy has provided information on the ground-state structure of the chromophore and its spectral properties when the protein environment is modified (van Thor et al., 1998; Bell et al., 2000). Some information was also obtained on the ground-state and first excited-state dynamics by time-resolved fluorescence spectroscopy (Chattoraj et al., 1996; Lossau et al., 1996; Didier et al., 2002; Winkler et al., 2002) and by resonance Raman spectroscopy (Esposito et al., 2001; Schellenberg et al., 2001; Tozzini and Nifosi, 2001). The excited-state decay kinetics is consistent with an excited-state proton transfer (EPST) from the chromophore to the protein (Chattoraj et al., 1996; Lossau et al., 1996; Wachter et al., 1997). Two deprotonated excited-state species of the chromophore, I^* and B^* , with similar lifetimes (3.3 ns and 2.8 ns, respectively) are thought to be mainly responsible for the GFP fluorescence (Striker et al., 1999).

One question still remains open: why do the denatured GFP and the chromophore isolated by enzymatic digestion of

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TABLE 1 Photophysical characteristics of the imidazolone derivatives in dioxane

Derivative	R ₁	R ₂	R ₃	Extinction coefficient (M ⁻¹ ·cm ⁻¹)	λ _{max} ^{abs} (nm)	λ _{max} ^{emis} (nm)	φ*
I-1	OH	Me	Me	25500	370	440	0.0001
I-2 [†]	OH	Me	(CH ₂) ₃ Me	26600	372	436	0.0001
I-3	OH	Ph	H	32600	399	467	0.0002
I-4 [†]	OH	3,4-diMeOPh	Me	28700	398	476	0.0002
I-5	H	Ph	H	13100	384	444	0.0012
I-6	H	Ph	OH	19700	385	443	0.0002
I-7	H	Ph	Ph	19500	380	453	0.001
I-8 [†]	H	4-MeOPh	H	26100	387	452	0.032
I-9 [†]	H	3,4-diMeOPh	H	23000	390	457	0.034
I-10 [†]	H	3,4-diMeOPh	Me	25900	386	470	0.045
I-11	MeO	Ph	H	15400	398	464	0.0002
I-12 [†]	MeO	Ph	Me	16800	394	471	0.0002
I-13 [†]	MeO	Ph	CH ₂ COOEt	25900	389	465	0.0002
I-14 [†]	MeO	Ph	CH ₂ Ph	20600	393	468	0.0002
I-15 [†]	MeO	4-NO ₂ Ph	H	22200	428	573	0.060
I-16 [†]	OCH ₂ COOEt	Ph	H	19100	397	461	0.0003
I-17	H, 2-MeO	Ph	H	26300	400	465	0.0022
I-18 [†]	N(Me) ₂	Me	H	39900	415	483	0.0009
I-19	N(Me) ₂	Ph	H	38800	454	520	0.0020
I-20 [†]	N(Me) ₂	4-MeOPh	H	37000	450	507	0.0023
I-21	N(Me) ₂	3,4-diMeOPh	H	50300	455	515	0.0025
I-22 [†]	CF ₃	Ph	H	22000	383	451	0.135
I-23 [†]	CN	Ph	H	24800	393	464	0.158
I-24 [†]	COOMe	Ph	H	28800	394	461	0.159
I-25 [†]	CN	4-MeOPh	H	29200	401	482	0.159
I-26 [†]	CN	3,4-diMeOPh	H	26500	405	482	0.189
I-27 [†]	COOMe	3,4-diMeOPh	H	32200	403	478	0.295
I-28 [†]	CN	4-NO ₂ Ph	H	20700	406	508	0.220
I-29 [†]	COOMe	4-NO ₂ Ph	H	25300	405	510	0.258

*Fluorescence quantum yields φ determined at 293 K. SEM of 0.001 for φ between 0.010 and 0.099 and of 0.005 for φ between 0.100 and 0.300.

[†]New products.

connectivities N₂-CA₂-CB₂-CG₂ and CA₂-CB₂-CG₂-CD₁, respectively. Energy mapping was conducted by 10° increment rotations of the imidazolone ring about its torsion angles. Rotations around the C₁-CA₁ bond were also considered but did not influence the potential energy mapping.

RESULTS

Photophysical properties of the imidazolone derivatives

The two main bands at ~395 and 475 nm in the absorption spectrum of wt-GFP are due to the neutral and anionic forms of the chromophore, respectively (Bell et al., 2000). Both forms are in equilibrium in the protein and correspond to the A and B forms in the four-state model elaborated by Weber et al. (1999) for the photophysical behavior of wt-GFP. The intense green fluorescence of the imidazolone chromophore at 508 nm seems to arise from the anionic B* form of the chromophore or from an intermediate I*. This last form results from the neutral A* after deprotonation by excited-state proton transfer from the imidazolone to the protein matrix (Lossau et al., 1996; Chatteraj et al., 1996). As the photophysical properties of the native chromophore are pH

sensitive, the electronic charge repartition in the molecule must play an important role in the fluorescence emission. Therefore, we synthesized several derivatives of the chromophore imidazolone with various electronic substituents and investigated the relationship between their electron distributions and photophysical properties.

In this study, the 4-hydroxybenzylidene-2,3-dimethyl-imidazolidinone (Fig. 1) was selected as the reference compound because it represents the smallest fluorescent moiety of GFP. The OH of the Tyr residue was conserved at atom CZ and the two aliphatic carbons of the peptidic backbone were replaced in CA₁ and CA₃ by methyl groups (compound I-1 called HBDI by Esposito et al., 2001). We substituted different radicals R₁, R₂, and R₃ in place of HO_y, CA₁, and CA₃, respectively (Table 1). First, derivatives of this molecule were synthesized with a phenyl as R₂ to mimic both the steric hindrance and the hydrophobic environment of part of the protein backbone. Then the hydroxyl in position R₁ was replaced by H to get free of the pH effect in this part of the molecule. Finally, various derivatives were obtained by substitutions in the positions R₁ and R₂ by electron-donating groups, such as MeO or N(Me)₂, or by electron-withdrawing groups, such as CF₃, NO₂, CN, or COOMe.

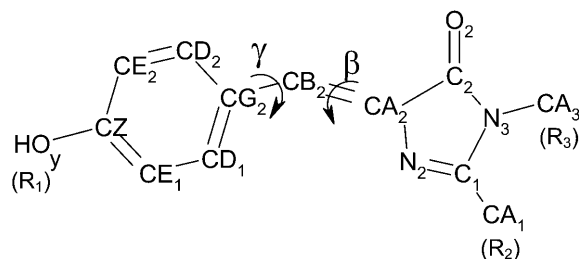


FIGURE 1 Model chromophore with atom labels according to the GFP crystal structure (Ormö et al., 1996). β and γ are the dihedral angles $N_2-CA_2-CB_2-CG_2$ and $CA_2-CB_2-CG_2-CD_1$, respectively.

The photophysical properties of these compounds (Table 1) were determined in dioxane, an aprotic solvent with a very small low-frequency dielectric constant ($\epsilon = 2.2$ at 25°C), commonly referred to as a good simulator of the hydrophobic environment of a chromophore buried inside a protein matrix. The absorbance and emission spectra of the most fluorescent compounds ($\phi > 0.1$) of Table 1 (I-22–I-29) are described here. Fig. 2 shows the spectra of I-23, I-27, and I-28, which are representative of the three kinds of spectra obtained and which are used to define three classes of derivatives [A], [B], and [C]. I-22 and I-24 show the same kind of spectra as I-23: all three correspond to compounds substituted by only one electron-withdrawing group in R_1 (Fig. 2 *a*). One notices that these compounds have structured absorption spectra, characterized by a main absorption peak with a shoulder on each side. They form class [A]. I-25 and I-26 show spectra similar to I-27 (Fig. 2 *b*). These compounds, substituted with an electron-acceptor in R_1 and a donor in R_2 , have much less structured spectra than the ones in Fig. 2 *a* and belong to class [B]. Fig. 2 *c* shows the spectra of I-28, which are very similar to the spectra of I-29. Both molecules have two electron acceptors in R_1 and R_2 and are red shifted in comparison with the previous ones. They are not structured at all and their peaks have greatly increased widths at half maximum. They correspond to class [C]. The broadening of the spectra in class [B] and [C] might be attributed to an increasing admixture of $\pi_{\text{ring}} \rightarrow A_{\text{acceptor}}^*$ or $D_{\text{donor}} \rightarrow \pi_{\text{ring}}^*$ states where π_{ring} and π_{ring}^* represent the bonding and antibonding orbitals of the aromatic rings respectively. In all cases, the S_1 - S_0 fluorescence spectrum is approximately the mirror image of the S_0 - S_1 absorption spectrum, which points to small overall differences between the nuclear configuration of the ground and excited states. However, the low intensity of the 0-0 band indicates considerably different equilibrium nuclear configurations in the ground and first excited states.

The extinction coefficient of the reference molecule I-1 at the absorption maximum is $25,500 \text{ M}^{-1} \text{ cm}^{-1}$ (Table 1) and there is a 105-nm blue shift in the absorption and a 70-nm blue shift in the emission compared to the anionic form of the chromophore in GFP. The substitution of R_2 by a phenyl ring (I-3) induces an increase in the extinction coefficient and

a marked red shift in both the absorption and emission spectra compared to I-1 ($\sim 30 \text{ nm}$) but the fluorescence quantum yield still remains negligible. The replacement of OH by H in compound I-5 induces a fall in the extinction coefficient with a small blue shift. The fluorescence quantum yield becomes nonnegligible when R_2 is substituted with the electron-donating groups 4-MeOPh or 3,4-diMeOPh in compounds I-8, I-9, and I-10.

In the literature, it appears that an electron-donating group in R_1 could be responsible for some fluorescence properties of the native anionic form of the chromophore (Bell et al., 2000). However, the substitution at R_1 by either MeO or $N(\text{Me})_2$, which are two strong electron-donating groups, does not have any influence on the fluorescence quantum yield that remains very low (I-11 to I-21). One can notice that a substitution in R_3 with a Me, CH_2COOEt or CH_2Ph does not significantly modify the spectroscopic characteristics. By contrast, a substitution at R_1 with electron-withdrawing groups induces a red shift, compared to the first compounds, both in the absorption maximum and in the emission maximum accompanied by an important increase in ϕ (from 0.158 to 0.220 for CN in R_1 and from 0.159 to 0.295 for COOMe in R_1) (I-23, I-25, I-26, I-28, and I-24, I-27, I-29, respectively). Among these fluorescent compounds, I-27 shows the greatest overall brightness, i.e., the product of the extinction coefficient and the quantum yield.

Table 1 clearly shows that the necessary and sufficient conditions for obtaining fluorescent imidazolones in dioxane with a quantum yield $\phi > 0.1$ at 20°C are: i), an electron-withdrawing group in R_1 and ii), a phenyl in R_2 , substituted with either an electron-donating group or an electron-withdrawing group. A large selection of fluorescent compounds is thus obtained with maximum absorption and emission wavelengths ranging from 383 to 406 nm and from 451 to 510 nm, respectively.

Fluorescence quantum yields and activation energies

After the excitation by an adequate radiation, and transition from the ground state S_0 to the first singlet excited state S_1 , the synthetic compounds may follow several theoretical deexcitation pathways. These can be described as one radiative process, the fluorescence, and several nonradiative processes. The nonradiative pathways include either a fast internal conversion or an intersystem crossing from the first electronic excited state S_1 to the T_1 triplet state. A third nonradiative pathway includes the excited-state intramolecular twisting about the two dihedral angles β ($N_2-CA_2-CB_2-CG_2$) and γ ($CA_2-CB_2-CG_2-CD_1$) followed by fast internal conversion after conical intersection. All these processes compete with each other.

In the following simplified kinetic scheme, two geometric forms of the molecules are considered, a planar *P* ($\beta = 0^\circ$, $\gamma = 0^\circ$) and a twisted *Tw* ($\beta \neq 0^\circ$ and/or $\gamma \neq 0^\circ$).

$P + h\nu \rightarrow P^*$	$S_0 \rightarrow S_1$	Excitation	
$P^* \rightarrow P + h\nu$	$S_1 \rightarrow S_0$	Fluorescence	(k_f)
$P^* \rightarrow P$	$S_1 \rightarrow S_0$	Radiationless deactivation via internal conversion and/or intersystem crossing	(k_{nr}^P)
$P^* \rightarrow T_w^* \rightarrow P$	$S_1 \rightarrow S_0$	Radiationless deactivation via conical intersection	(k_{nr}^{Tw})

The fluorescence quantum yield ϕ of a fluorescent compound is thus given by:

$$\phi = \frac{k_f}{k_f + k_{nr}^P + k_{nr}^{Tw}} \quad (1)$$

Low-temperature behavior

Some low-temperature experiments were undertaken with the protic solvent EPA used for glass-forming solutions (see Materials and Methods). Temperature was slowly decreased from 293 K to 83 K in a thermostated cuvette in the fluorescence spectrophotometer, and steady-state fluorescence spectra were recorded every 15 min. The variation of the fluorescence quantum yield ϕ as a function of temperature is shown in Fig. 3 for three derivatives having different fluorescent properties at 293 K: I-1, which does not present any fluorescence at this temperature in water, EPA or

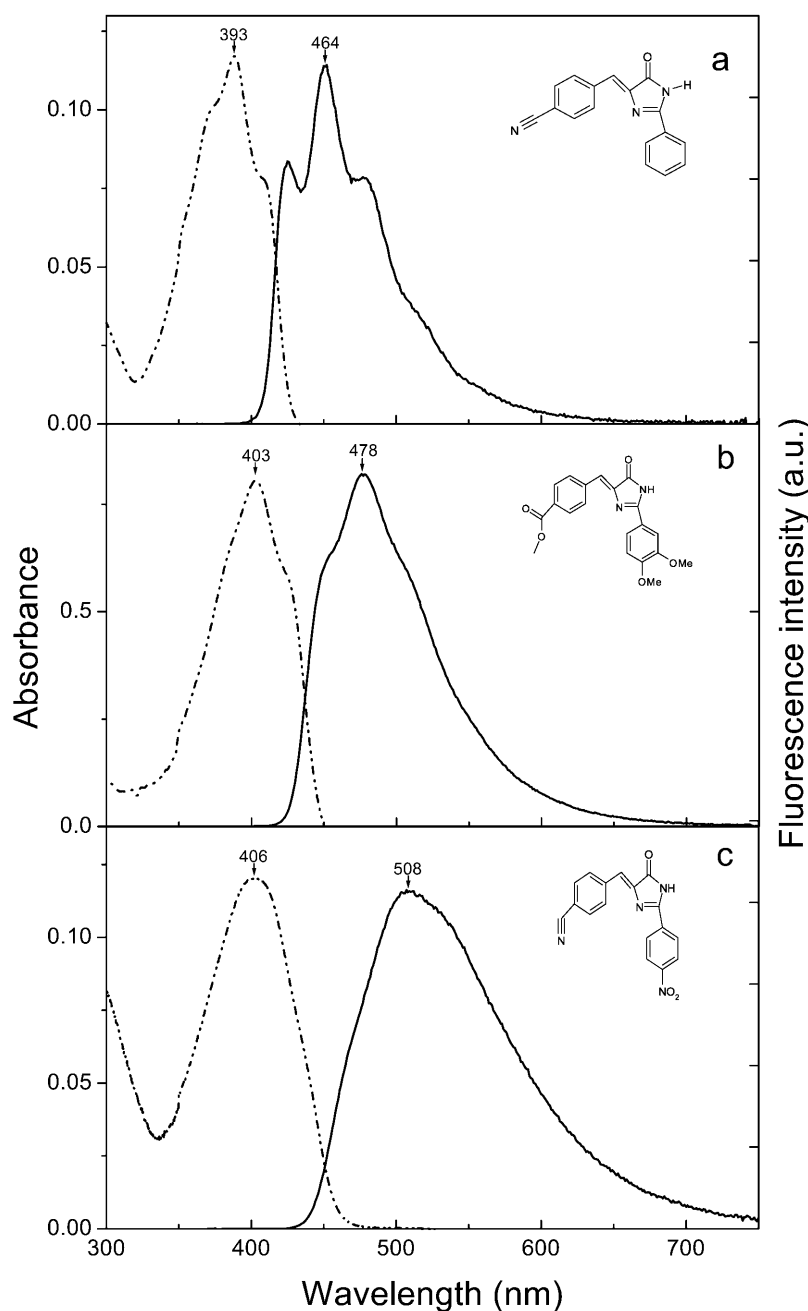


FIGURE 2 Room-temperature steady-state absorption spectra (*dashed*) and fluorescence emission spectra (*solid*) of GFP derivatives in dioxane: I-23 (*a*); I-27 (*b*); and I-28 (*c*). The absorption and fluorescence spectra were normalized to approximately the same maximum amplitude.

dioxane, and I-24 and I-27, which are fluorescent in EPA and dioxane.

The fluorescence of I-1 is undetectable until 133 K, a temperature near the glass-forming temperature of the solvent. The fluorescence intensity strongly increases with lower temperatures without reaching a maximum. This can be explained by the progressive rigidification of I-1 in the cooled, solution, which induces an important decrease in the nonradiative deexcitation processes. For I-27 and I-24, the fluorescence intensity increases progressively with decreasing temperatures and reaches a constant level corresponding to maximal fluorescence quantum yields of 0.42 and 0.78, respectively. It is noticeable that at 80 K the fluorescence quantum yield of I-24 reaches the high value of 0.79 observed at room temperature for the wt-GFP chromophore in the protein.

At low temperature, the probability of the molecules to be in a twisted conformation is very low in comparison to the probability of the planar conformation. Consequently, k_{nr}^{Tw} in Eq. 1 becomes negligible. Because the maximal quantum yield ϕ_0 obtained at 80 K corresponds to:

$$\phi_0 = \frac{k_f}{k_f + k_{nr}^P} \quad (2)$$

the variation of the rate constant k_{nr}^{Tw} with temperature can be calculated by means of the general formula described by Saltiel and Sun (1989) in the case of transstilbene twisting.

$$k_{nr}^{Tw} = k_f \left(\frac{1}{\phi} - \frac{1}{\phi_0} \right) = \alpha \times \exp^{-E_a/k_B T} \quad (3)$$

where α and E_a are the frequency factor and the activation energy of the first excited-state S_1 . T represents the absolute temperature and k_B the Boltzmann constant.

The radiative decay rate constant, k_f , equal to the ϕ/τ ratio, was assumed to be equal to $1.4 \times 10^8 \text{ s}^{-1}$, taking the fluorescence quantum yield of I-27 (Table 1) and its fluorescence lifetime determined for other purposes as $2.1 \times 10^{-9} \text{ s}$. This value of k_f was assumed to be temperature and solvent independent, and used for all the compounds based on their structural similarity.

The curves displayed in the inset of Fig. 3 are the corresponding Arrhenius plots for the k_{nr}^{Tw} rate constant. However, because the viscosity of the medium increases with decreasing temperatures and, consequently, the activation energy barrier is medium enhanced (Saltiel and Sun, 1989), it is interesting to note the non-Arrhenius behavior and the “characteristic knee” already mentioned by Litvinenko et al. (2003) and attributed to the onset of strong caging at the critical crossover temperature T_c of the glass-forming solvent. In our case, because the solvent is a mixture of three components, the T_c value is not known but the “knee” appears well above the glass transition temperatures T_g of ethanol (95 K) and isopentane (65 K) (Carpentier et al., 1967; Sugisaki et al., 1968) and therefore might well correspond to the T_c value of EPA.

Relationship between activation energies E_a and fluorescence quantum yields

The main nonradiative deexcitation processes from the singlet excited state evoked so far in the literature are a *cis-trans* photoisomerization of the exo-methylene double bond (Niwa et al., 1996; Weber et al., 1999) or more limited rotations about the two dihedral angles β and γ (Phillips, 1997; Kummer et al., 1998; Voityuk et al., 1998; Chen et al., 2001; Webber et al., 2001). Therefore, the theoretical activation energies of the different conformations obtained

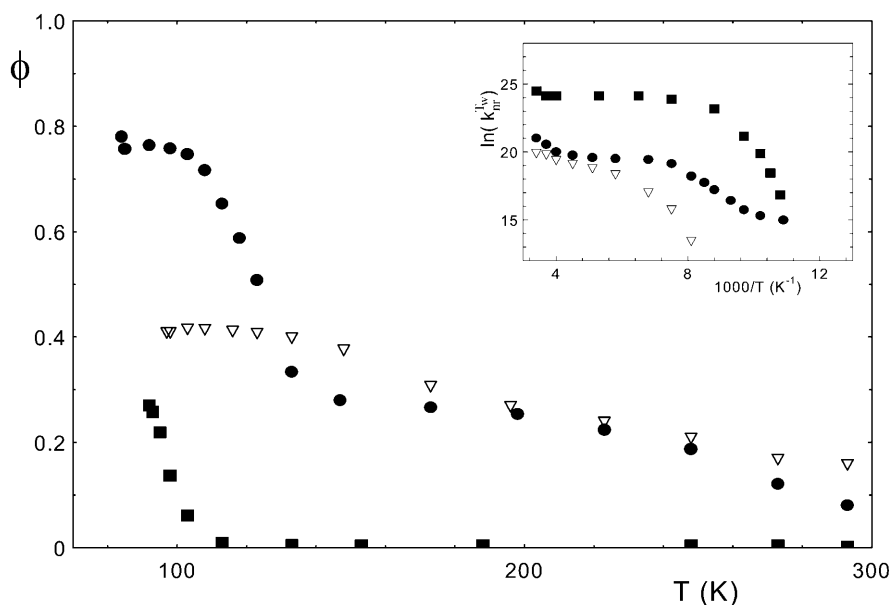


FIGURE 3 Low-temperature effect on the fluorescence quantum yields of I-1 (■), I-24 (●), and I-27 (▽). The maximal SEM on the ϕ values are of 0.005. The inset shows the corresponding Arrhenius plots for the k_{nr}^{Tw} rate constant. Measurements were performed in EPA. For details, see Materials and Methods.

after internal rotation about these angles were determined for some 20 compounds by semiempirical calculations with the MOPAC 93 program. The energies of I-1 and I-23 were also calculated with the *ab initio* Gaussian 98 program (CIS method, STO-3G basis set) (Frisch et al., 1998) and a good correlation was observed (data not shown).

The energy mapping was conducted for the ground and first excited states by simultaneous 10° increments of both dihedral angles. The potential energy surfaces (PES) were identical for all rotations of the R₂ substituent around the C₁-CA₁ σ bond, which could be explained by the axial symmetry of nearly all the substituents. This rotation angle was further kept at 0° to give a planar conformation to the imidazolidinone ring coupled to R₂. As an example of the results, the PES of I-23 are represented in Fig. 4 *a*. For each compound, we checked that the 0-0 transition energies estimated from the intersection of the excitation and emission spectra, are reasonably well predicted by the MOPAC calculations: the differences amount at most to 3%. For β = 90° an “avoided crossing” between the S₀ and the S₁ PES is observed for all the derivatives (Voityuk et al., 1998).

It is possible to restrict the two-dimensional energy diagram of I-23 to an energy profile because i), the energy surfaces present a symmetry around β = 90° and ii), the lowest energies of the excited state are obtained with γ = 0°. The profile plotted in Fig. 4 *b* represents the relative variation of the first excited-state energy $E_{S_1}(\beta)$ during the rotational motion of the imidazolidinone ring around β, compared to the energy $E_{S_1}(0)$ of the planar conformation of the molecule with β = γ = 0°. The same type of profile was obtained for all the derivatives. The maximum of E_{S_1} was always observed for β between 40° and 50°. These maxima, given in Table 2, represent the activation energy, E_a , of the barriers that must be overcome to reach the avoided crossing region, i.e., the nonradiative decay channel. They seem to play a key role in the photophysical properties of the molecules because their height is linked to the fluorescence quantum yield.

The relationship between the activation energies E_a of 12 compounds and their fluorescence quantum yields at room temperature in dioxane (Table 1) was investigated.

From Eqs. 1 and 3, one obtains:

$$\frac{1}{\phi} = 1 + \frac{k_{nr}^P}{k_f} + \frac{\alpha}{k_f} \exp(-E_a/k_B T) \quad (4)$$

Assuming that the k_{nr}^P/k_f ratio in Eq. 1 does not significantly differ at room temperature from one compound to another because all the compounds show the same basic chemical structure, the semilogarithmic variation of $(1/\phi - 1 - k_{nr}^P/k_f)$ was plotted as a function of $E_a/k_B T$ (Fig. 5). The constant ratios were solved to fit the experimental data linearly with a slope of -1. The values obtained for k_{nr}^P/k_f (2.39) and α/k_f (1470) are consistent

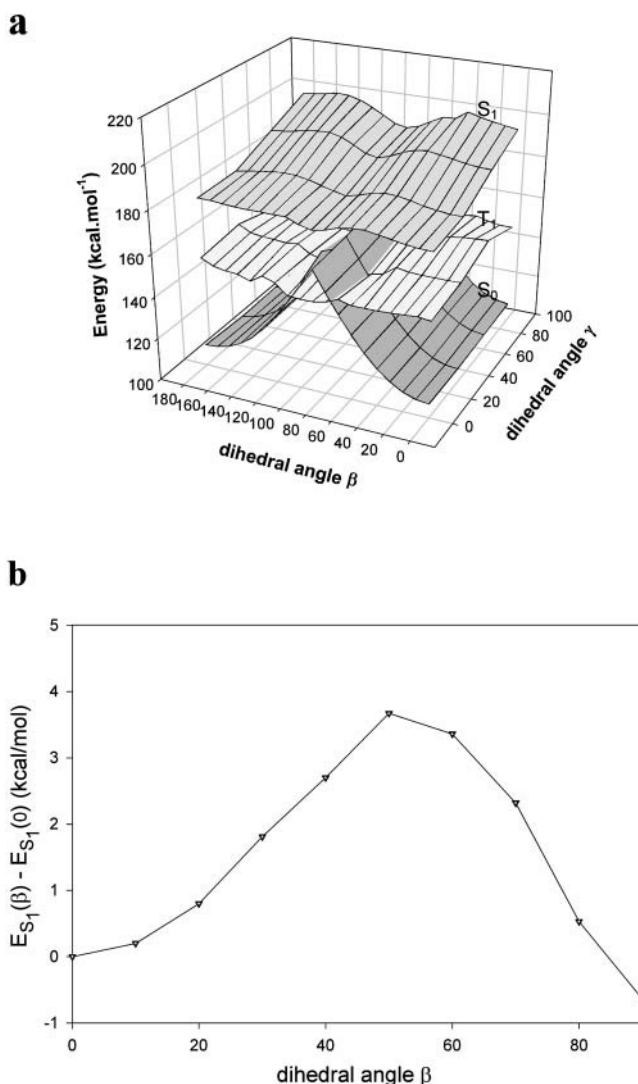


FIGURE 4 Energy diagrams of I-23: (a) Calculated two-dimensional potential energy surfaces for the S₀, S₁, and T₁ states as a function of the dihedral angles β and γ; (b) Energy profile of the S₁ state as a function of the dihedral angle β (from 0° to 90°), with γ held fixed at 0°. The zero value corresponds to the energy of the molecule with (β, γ) equal to (0, 0). Calculations were made with the MOPAC program.

with the proposal that the avoided crossing between the ground and first excited state is the main nonradiative channel for these molecules because $\alpha \gg k_{nr}^P$. The points corresponding to I-1, I-3, and I-5 that are clearly outside the fitting curve should be fitted by another curve parallel to the first one with a higher Y-intercept, reflecting a higher α/k_f ratio. This variation can be explained by comparing the chemical structures of the different compounds. Actually, I-1, I-3, and I-5 are the simplest derivatives that do not possess any electron-donating or electron-withdrawing substituents in R₁ and R₂ and thus can freely twist around the exo-methylene bond.

TABLE 2 Theoretical maximum activation energies of various compounds

Derivative*	R ₁	R ₂	Energy (kcal mol ⁻¹)
I-1	OH	Me	1.96
I-3	OH	Ph	2.21
I-5	H	Ph	2.53
I-8	H	4-MeOPh	1.87
I-22	CF ₃	Ph	3.27
I-23	CN	Ph	3.67
I-25	CN	4-MeOPh	4.26
I-24	COOMe	Ph	3.42
I-26	CN	3,4-diMeOPh	4.30
I-27	COOMe	3,4-diMeOPh	7.14
I-28	CN	4-NO ₂ Ph	3.68
I-29	COOMe	4-NO ₂ Ph	5.81

*In all cases, R₃ is H

Effect of extrinsic factors on the fluorescence quantum yields

A total rigidification of I-24 leads to the fluorescence yield of 0.79 seen in the wt-GFP chromophore. However, when the chromophore is isolated from the proteic cage, the fluorescence quantum yield falls down. It was thus interesting to study the influence of the environment on the fluorescence quantum yields of the synthesized derivatives by using three solvents of various hydrophobicity, viscosity, or polarity.

As dioxane is miscible in all proportions with highly polar solvents like water, the fluorescence emission of the various derivatives was measured in 1/1000 (v/v) dioxane/water solutions. The spectra obtained (data not shown) are much less structured in the aqueous medium than in pure dioxane and the fluorescence emission maximum shifts to the red by 50–100 nm, possibly as a result of the increase in solvent

polarity. However, the fluorescence quantum yields of the compounds never exceed 0.003 (Table 3). One deduces that in water the radiationless relaxation via internal conversion is largely preponderant over the radiative processes. Because all the molecules investigated here possess a large number of lone pair orbitals and hence are good hydrogen-bond acceptors, a quenching process through hydrogen bonding with solvent molecules must occur (see Biczok et al., 1997 and references given therein). To strengthen this conclusion we studied the effect of adding traces of water to a solution of I-27 in dioxane whose quantum yield is 0.295. Adding water, even at a level as low as 1.4% (v/v) to the solution immediately results in a drop of ~30% in quantum yield.

The fluorescence quantum yields were also determined in 100% glycerol solutions for all the derivatives with a fluorescence quantum yield in dioxane $\phi_{\text{diox}} > 0.10$ (Table 3). In all cases, we observe a 15- to 100-fold increase in the ϕ_{gly} values in this highly viscous medium relative to the results obtained in water. However, the values are always lower than those obtained in the hydrophobic low polarity solvent dioxane. The ratio $\phi_{\text{diox}}/\phi_{\text{gly}}$ leads again to the same three classes of derivatives distinguished by their photo-physical characteristics in dioxane and that depend mainly on the R₂ substituent. The number of hydrogen bonds between the H-bond acceptor atoms of the substituents in R₁ and R₂ and the solvent as well as the solvent-solute dipolar interactions might explain the differences between the three classes. It appears that hydrogen bonding would have an opposite effect to that induced by the viscosity increase and favor the nonradiative process. In class [A], as both the ϕ_{diox} and ϕ_{gly} are similar, one can conclude that through a combined viscosity increase and H-bond number decrease from water to glycerol, it is nearly possible to reach the same fluorescence quantum yield as in dioxane. In classes [B] and

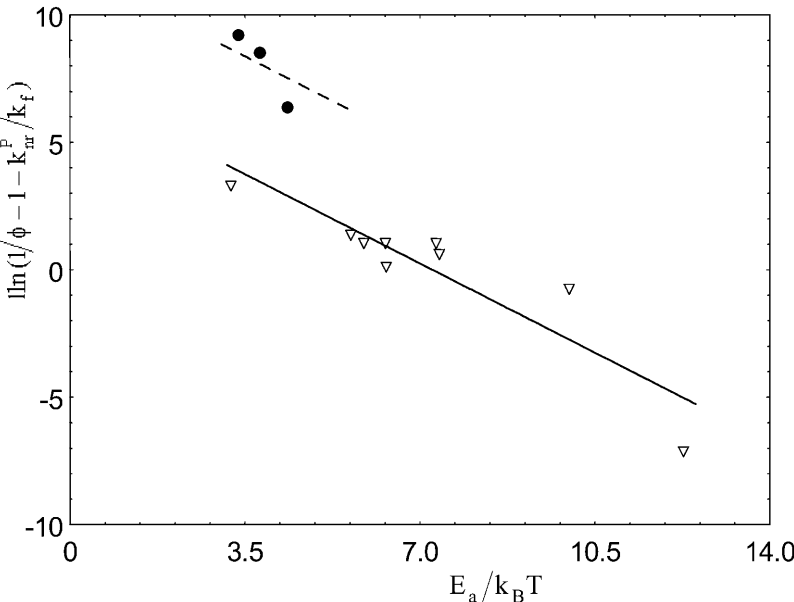


FIGURE 5 Semilogarithmic variation of $(1/\phi - 1 - k_{\text{tr}}^{\text{P}}/k_{\text{f}})$ as a function of $E_{\text{a}}/k_{\text{B}}T$. Derivatives with $\phi \geq 0.1$ (∇) and $\phi < 0.1$ (\bullet). The fluorescence quantum yields were determined in dioxane at 293 K.

TABLE 3 Fluorescence quantum yields in water, glycerol, and dioxane

Class	Derivative*	R ₁	R ₂	ϕ_{water} ($\epsilon^{\dagger} = 80$)	ϕ_{glycerol} ($\epsilon = 40$)	ϕ_{dioxane} ($\epsilon = 2$)	$\phi_{\text{diox}}/\phi_{\text{gly}}^{\ddagger}$
[A]	I-22	CF ₃	Ph	0.003	0.118	0.135	1.14
	I-23	CN	Ph	0.003	0.139	0.158	1.14
	I-24	COOMe	Ph	0.0016	0.156	0.159	1.02
[B]	I-26	CN	3,4-diMeOPh	0.0006	0.069	0.189	2.74
	I-27	COOMe	3,4-diMeOPh	0.0015	0.100	0.295	2.95
[C]	I-28	CN	4-NO ₂ Ph	0.0015	0.023	0.220	9.56
	I-29	COOMe	4-NO ₂ Ph	0.0006	0.014	0.258	18.4

*In all cases, R₃ is H.[†] ϵ is the low-frequency dielectric constant of the solvent.[‡]Ratio of the fluorescence quantum yields obtained in dioxane and glycerol.

[C], the ϕ_{gly} are lower than the values obtained in class [A], although I-26 to I-29 are the most fluorescent compounds in dioxane. The more numerous H-bonds between their bulky substituents in R₂ (particularly on the 3,4-diMeOPh) and the solvent would limit the fluorescence quantum yield increase. The differences observed between classes [B] and [C] could only be attributed to the opposite electronic properties of the substituents, 4-NO₂Ph and 3,4-diMeOPh.

In view of the high fluorescence quantum yield of wt-GFP, free dissolved O₂ does not seem to enter into the GFP protein matrix once the β -barrel is formed (Ormö et al., 1996). The chromophore is only in contact with the oxygen atom of structural water molecules through hydrogen bonds. But in the present work, the derivatives are surrounded by the solvent, which may dissolve molecular O₂. To check a possible quenching effect of O₂ on their fluorescence, some experiments were performed on samples deaerated by bubbling argon. The fluorescence quantum yields of I-27 obtained with and without dissolved O₂ are compared (Table 4). This derivative was chosen because it is the most fluorescent compound in both water and dioxane, so that even small quantum yield variations are detectable. Table 4 shows that dissolved O₂ has no significant inhibitory effect in EPA or dioxane. In water, the quantum yield slightly increases in the absence of O₂ but not significantly, so that dissolved O₂ cannot explain the loss of the fluorescence properties of the isolated chromophore.

DISCUSSION

In wt-GFP, the fluorescence spectrum is characterized by a strong emission peaking at 508 nm, which seems to

originate from two deprotonated states, an environmentally unrelaxed intermediate I* state or a relaxed anionic state B* (Creemers et al., 1999; Chattoraj et al., 1996). Its high quantum yield ($\phi = 0.79$) results from a conjunction of several factors. In the crystal structure of wt-GFP (Protein data bank access code 1ema), it appears that the chromophore is highly protected from the bulk solvent. Indeed, it was shown that the viscosity of the surrounding medium affected the fluorescence decay of the isolated chromophore (Kummer et al., 2002) but not that of the GFP (Suhling et al., 2002a), which is however sensitive to the refractive index (Suhling et al., 2002b). In the protein, the chromophore is surrounded by both apolar and polar residues, including a number of charged residues (Ormö et al., 1996). On one side of the chromophore, there is a large cavity containing only water molecules that contribute to a hydrogen bonding network linking the buried side chains of Glu-222 and Gln-69 (Chen et al., 2001). On the opposite side, there are polar interactions through hydrogen bonds between the phenolic hydroxyl of the chromophore and His-148, Thr-203, and Ser-205. There are also hydrogen bonds bridging the carbonyl of the imidazole ring to Glu-94 and Arg-96; this last residue is in a protonated form, which contributes to the partial negative charge residing on this carbonyl oxygen. All the surrounding residues and this dense hydrogen bonding network contribute to stabilize the chromophore in a nearly planar conformation.

The absence of fluorescence of the isolated chromophore has already given rise to studies on some models of the chromophore (Kojima et al., 1998; Webber et al., 2001; Kummer et al., 2002). Here we studied a large selection of synthetic derivatives of the imidazolidinone with various substituents, introduced both to increase the rigidity of the molecule and to modify the electron repartition in the molecule. Theoretical results with MOPAC 93 show that a crucial parameter as regards fluorescence ability is the activation barrier for a twist around the β angle. In non- or little-substituted compounds, the molecule is free to twist around the β and γ angles. It therefore emits very little fluorescence but undergoes a fast radiationless decay by internal conversion after reaching an avoided crossing. If the

TABLE 4 Influence of molecular oxygen on the fluorescence quantum yield of I-27 at 20°C

Solvent	Water	Glycerol	EPA	Dioxane
Dielectric constant	80	40	24	2
ϕ (oxygenated solution)	0.0020	0.100	0.159	0.295
ϕ (without O ₂)	0.0027	nd*	0.159	0.285

*Not determined in 100% glycerol because of argon bubbles.

TABLE 5 Photophysical characteristics of some mutants of GFP

Mutations	Extinction coefficient (M ⁻¹ cm ⁻¹)	$\lambda_{\text{max}}^{\text{abs}}$ (nm)	$\lambda_{\text{max}}^{\text{emis}}$ (nm)	ϕ
None*	25500	396, 475	508	0.79
I167T [†]	—	396, 471	502	—
N146I, M153T, V163A*	—	452	505	—
F99S, M153T, V163A (cycle3) [‡]	30000	397	506	0.79
F64L, I167T, K238N [§]	11500	395, 468	506	0.84
E222Q [¶]	—	476	506	0.5
F64L [§]	18300	396, 473	508	0.85
E222G	—	475	510	—
T203I	—	399	511	—
S202F, T203I*	20000	399	511	0.60
T203I, S72A, Y145F*	29000	399	511	0.64

*Tsien (1998).

[†]Heim et al. (1994).[‡]Patterson et al. (1997).[§]Palm et al. (1997).[¶]Jung et al. (2000).^{||}Ehrig et al. (1995).

molecule is substituted with certain well-defined bulky substituents, the equilibrium between both radiative and nonradiative deexcitation pathways is modified and the compounds are fluorescent. The fluorescence quantum yields of our compounds were linked to the activation energies by introducing a new term corresponding to their twisted form in the classical scheme used for rigid aromatic hydrocarbons (Birks, 1973). Reducing or blocking the rotation about the exo-methylene double bond by either increasing the viscosity of the solvent or by lowering the temperature leads to increased fluorescence quantum yields in all the derivatives tested. Moreover a correlation is found between the inductive effect of the substituents and the quantum yield. It appears that all compounds with a high quantum yield (I-22 to I-29) possess an electron-withdrawing substituent in R₁ and a phenyl group in R₂ either not substituted or substituted by an electron-withdrawing or an electron-attracting substituent.

Additional information is obtained from the GFP mutants. Mutations of the residues of the β -barrel have various consequences, depending on whether they are within the chromophore or just in its vicinity in the folded protein. When mutations are located within the chromophore (mutations of residues Ser-65, Tyr-66, or Gly-67), the structural formula of the chromophore is completely modified (Cubitt et al., 1995; Wachter et al., 1997; Tsien, 1998). These mutants will not be discussed, because it is obvious that the spectral characteristics will be affected. More interesting are the mutations in the vicinity of the chromophore that indirectly induce i), a modification in the hydrogen bonding network and ii), steric and electrostatic

changes in its close environment. These changes induce modifications in the absorption and emission spectra. Table 5 reports the photophysical characteristics of some of these mutants that are fluorescent. The $\lambda_{\text{max}}^{\text{abs}}$ indicate that some mutants exist in the neutral and anionic forms, like wt-GFP. Others are mainly in the anionic form (E222G) or the neutral form (cycle 3, T203I). They are all emitting with a $\lambda_{\text{max}}^{\text{emis}}$ between 502 and 511 nm, which indicates that it is always the anionic form that is fluorescent, probably after a fast excited-state proton transfer from the neutral form to the surrounding protein matrix, as in wt-GFP (Chattoraj et al., 1996; Lossau et al., 1996; Wachter et al., 1997). A comparison with our results in Table 1 leads to the conclusion that compounds I-28 and I-29 (class [C] compounds) show similar spectroscopic characteristics. One could put forward the hypothesis that the rigidity of I-28 and I-29 and their electron distribution are similar to that found in the chromophore of these mutants, i.e., an electron attractive effect of the protein surroundings near the R₁ and R₂ positions.

Because the chromophore in GFP is linked to its environmental amino acids through hydrogen bonds and is under the electrostatic influence of its neighboring residues, our strategy to obtain further insight in GFP photophysics will be to bury our fluorescent products in hydrophobic protein pockets. Moreover thanks to the dependence of their emission on the presence of hydrogen-bond donors, on viscosity, and on rotational mobility, these synthetic derivatives may also find use as fluoroprobes of the microenvironment of proteins and other biological macromolecules.

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